

The Un-Natural PNGB Higgs

Luca Vecchi

(University of Maryland)

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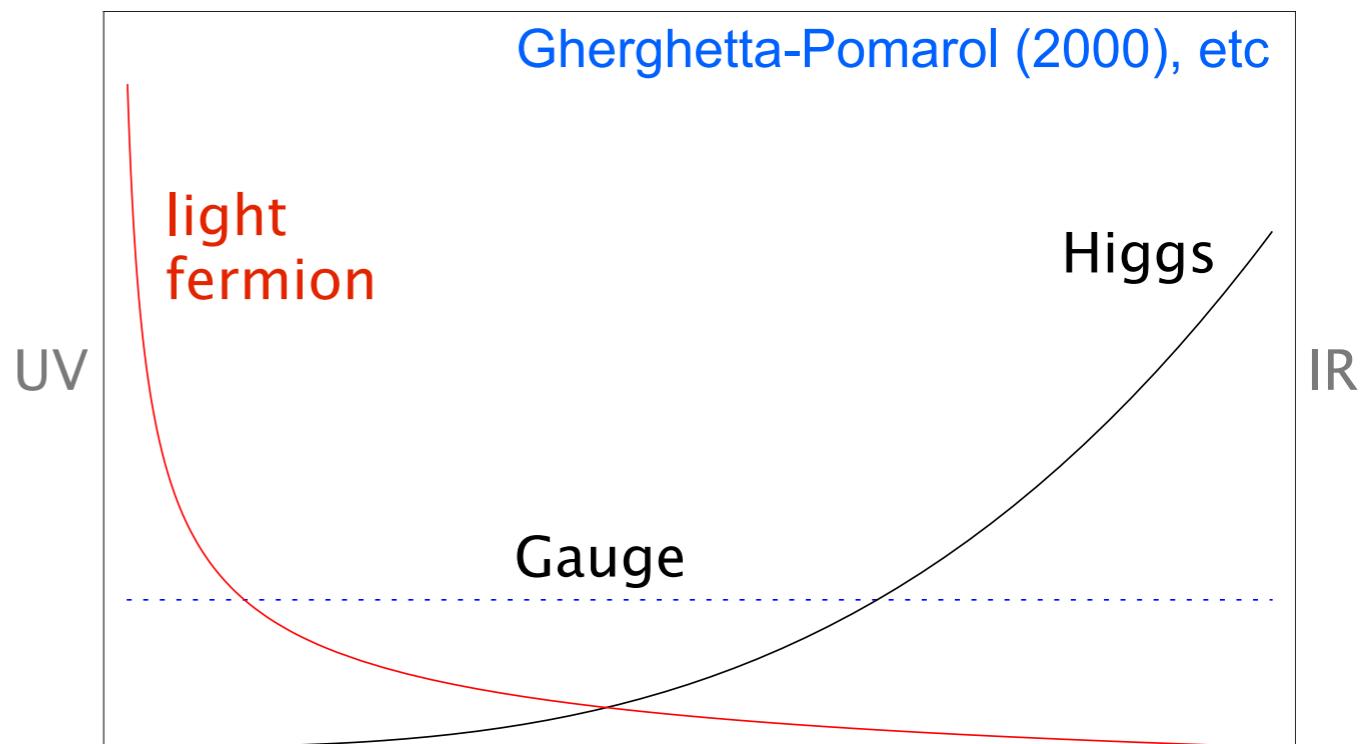
PNGB Higgs models are still a viable alternative to SUSY

a problem then ————— Flavor ————— a virtue now

$$\frac{\bar{q}q\mathcal{O}}{\Lambda^d}$$

$$\epsilon \bar{q}\mathcal{O} \longrightarrow y \propto \epsilon_L \epsilon_R$$

hierarchy
via RG



Outline:

- * Status after Higgs discovery (see Matthias talk)
 - mass at 125 GeV
- * Natural Vs “Un-natural” PNGB Higgs:
 - **Why un-naturalness?**
- * Signatures of un-naturalness

The pseudo-NGB Higgs

Georgi-Kaplan ('80s)

Broad picture: NDA

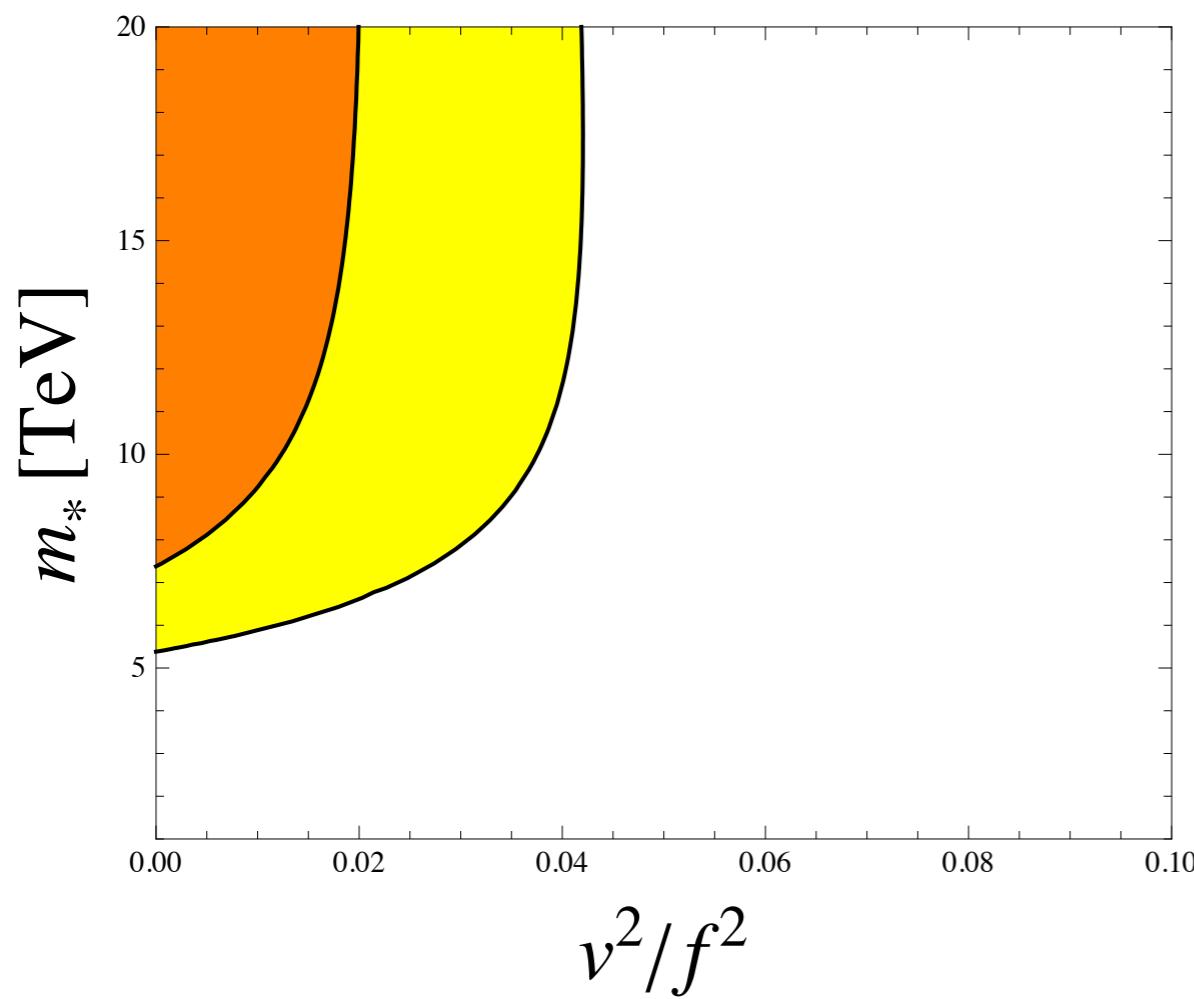
Georgi et al.
Contino et al. (2006)
Giudice et al. (2007)

$$m_* \sim g_* f \quad \text{one coupling, one mass scale}$$

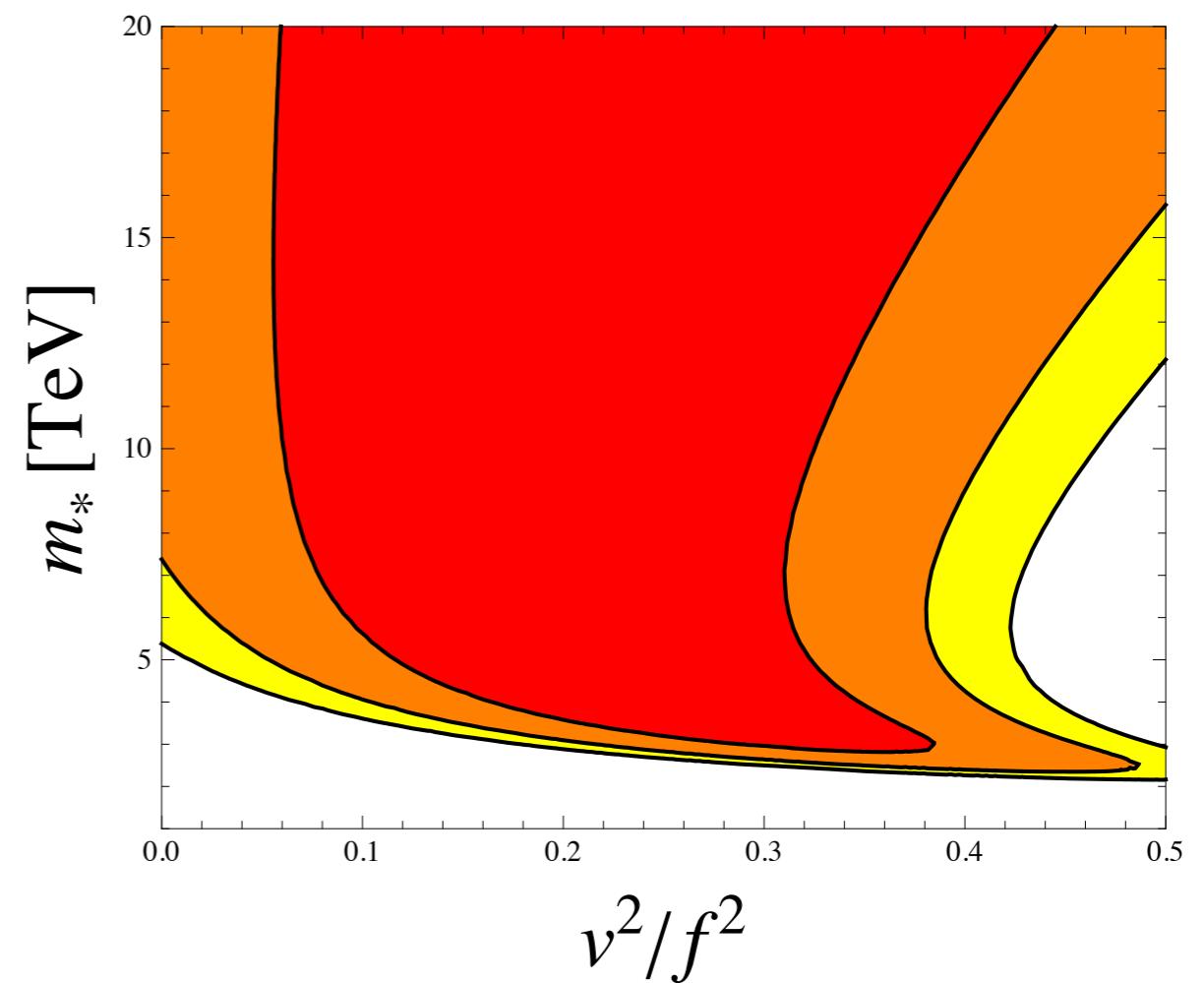
EW precision data

GFitter (2014) U=0

$$\Delta T = \Delta T_h$$



$$\Delta T = \Delta T_h + \frac{v^2}{f^2}$$



conservative

vs

optimistic ($Z \rightarrow b\bar{b}$?)

Higgs mass

$$V = \frac{y_t^2 N_c}{16\pi^2} m_*^2 f^2 \left[a \sin^2 \frac{h}{f} + b \sin^4 \frac{h}{f} + \dots \right] \text{ plus gauge and/or exotic contributions}$$

$$\frac{v^2}{f^2} \sim \frac{a}{b} < 1$$

$$m_h = 100 \text{ GeV} \sqrt{b \frac{g_*^2}{y_t^2}}$$

$$m_h = 100 \text{ GeV} \sqrt{b \frac{g_*^2}{y_t^2}}$$

$g_* = ?, b \ll 1$

larger tuning

$g_T \ll g_*, b \sim 1$

light top partners

Ex: minimal SO(5)/SO(4) with bulk fermions in the irreps 4,5,10...

Pomarol et al. (2012)

Panico et al. (2012)

Marzocca et al. (2012)

...

tuned, but
potentially generic

vs

split spectrum

Example:

$SU(4)/Sp(4)$ with $\epsilon \bar{q} \mathcal{O}$, and $\mathcal{O} \sim \mathbf{4} \in SU(4)$

$$V \propto \sin^2 + \frac{y_t^2}{g_*^2} \sin^4 + \dots \Rightarrow \begin{cases} b \sim \frac{y_t^2}{g_*^2} \\ \lambda = c \frac{N_c y_t^4}{4\pi^2} \ln \frac{m_*}{m_h} \end{cases}$$

Lesson:

Once tuning ($v/f \ll 1$) is swallowed $\sim \frac{m_*^2}{m_t^2}$

- 1) split spectrum is not necessary
- 2) coupling strength is not constrained

Flavor

**typical spectrum >10 TeV
with flavor anarchy**

Bauer et al. (2010)

Keren-Zur et al. (2013)

...

light top partners

$m_T > 10$ TeV ($\Delta F=2$) \iff
tuning or symmetries

Keren-Zur et al. (2013)

Barbieri et al. (2013)

...

hierarchy from
flavor anarchy

vs

flavor symmetries
needed

Natural

Vs

Un-Natural

Tuning

<1/10

<1/1000

EW data

with luck

generically

Flavor

non-generic

very attractive

Higgs mass

split spectrum

generic spectrum

Natural

Vs

Un-Natural

Tuning

<1/10

EW data

with luck

Flavor

non-generic

Higgs mass

split spectrum

Non-generic

**natural but
non-generic**

?

OR

Tuned
<1/1000

generically

very attractive

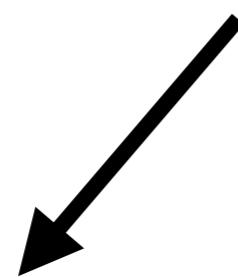
generic spectrum

**generic and
fine-tuned**

How much tuning? and “Why”?

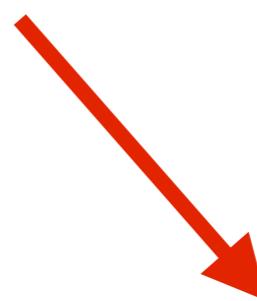
WIMP “miracle”?

$$\sigma v = \frac{g_*^4}{4\pi m_*^2} \Rightarrow m_* \sim (3 - 5)g_*^2 \text{ TeV}$$



light composites
annihilate into SM

mass in the 1s of TeV



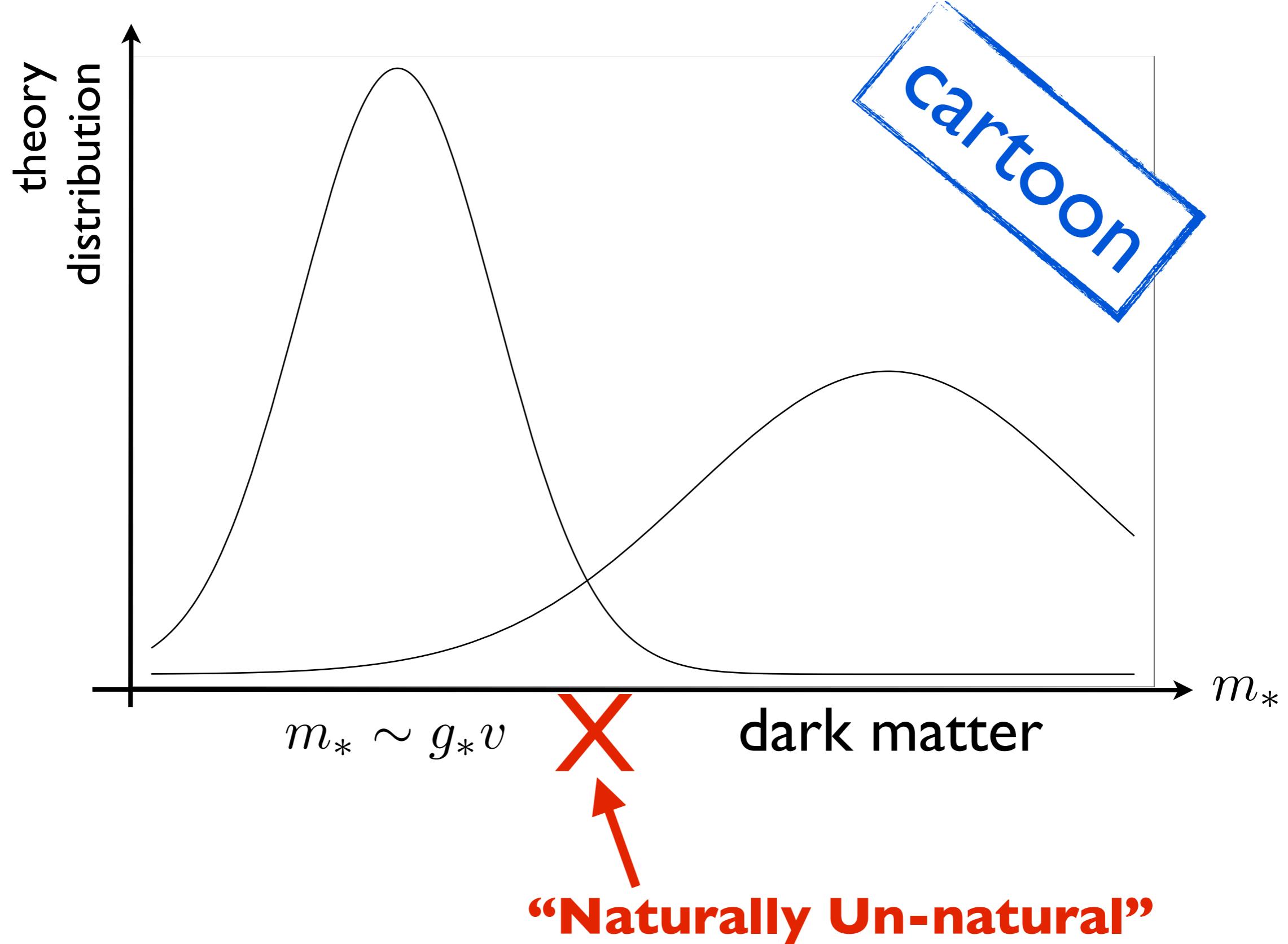
typical composites
annihilate into BSM

mass in the 10s of TeV



tuning is generic
("natural" in a broader sense)

One parameter, Two observable to explain “naturally”...



WIMPs in warped extra dimensions?

Baryon DM Nussinov (1985) (soliton in 5D)

pp-->pi's has $\sigma_{p\bar{p}}v \sim \text{few} \times 10 \text{ mb}$

$$\sigma_{X\bar{X}}v \sim \sigma_{p\bar{p}}v \left(\frac{\text{GeV}}{m_X} \right)^2 \sim \text{few pb} \left(\frac{100 \text{ TeV}}{m_X} \right)^2 \text{“upper bound”}$$

Heavy WIMPs in warped extra dimensions are **generic**

Stability ✓

B&L is typically enough to have stable particles (SUSY)

$$(-)^{B+L+2s}$$

Ex: Majorana fermion, boson with non-even B, etc.

in warped 5D B&L are spontaneously broken gauge symmetries.

Some “hadron” carries B&L: $\lambda \bar{q} \mathcal{O}$

Annihilation into BSM ✓

$$m_* \sim (3 - 5)g_*^2 \text{ TeV}$$

Majorana DM

$U(I)_X$ broken in the bulk by 5D Majorana mass (i.e. no chiral symmetry)
 \implies heavy DM

XX-->KK-fermions, Goldberger-Wise, etc. dominates

Baryonic (Leptonic) DM

Ex: scalar with non-even B (L)
 \implies **XX-->U(I)_B gauge fields is generic**

A testable picture: Signatures of Un-Naturalness

Dark matter
Colliders

...

WIMP detection

Indirect detection and colliders: model-dependent $q^2 \sim m_X^2$

Direct detection: model-independent $q^2 \ll m_X^2$

“heavy” WIMPs \Rightarrow determine only σ/m

no “standard pNGB model” \Rightarrow EFT approach [Bagnasco et al. \(1994\), LV \(2013\)](#)

$$\frac{d\sigma}{dE_R} \sim \left(\frac{\mu}{m_*}\right)^{2d-4} \frac{1}{m_*^2 E_R}$$

SM operator!

O_α^A	$g' B^{\mu\nu}$	$H^\dagger i \overleftrightarrow{D}^\mu H, \bar{q} \bar{\sigma}^\mu q$	$g_s^2 GG, \bar{q} H q, H^\dagger H$	$g_s^2 G_\alpha^{\{\mu} G^{\nu\}}{}^\alpha, \bar{q} \bar{\sigma}^{\{\mu} i D^{\nu\}} q$
$\mathcal{T}^A(\mathbb{C})$	$\varepsilon^{\mu\nu\alpha\beta} \mathbf{v}_\alpha \mathbf{s}_\beta, \mathbf{v}^\mu \mathbf{s}^\nu$	\mathbf{v}^μ	$\mathbf{1}$	$\mathbf{v}^\mu \mathbf{v}^\nu, \mathbf{s}^\mu \mathbf{s}^\nu, \mathbf{s}^\mu \mathbf{v}^\nu$
$\mathcal{M}_N _{\text{SI}}$	$\frac{(\vec{v}_\perp \wedge \vec{s}) \cdot \vec{i}q}{\vec{q}^2}, \frac{\vec{s} \cdot \vec{i}q}{\vec{q}^2}$	1	1	$1, 1, \vec{s} \cdot \vec{v}_\perp$
$\mathcal{T}^A(\mathbb{R})$	$\mathbf{s}^\mu \mathbf{i} \mathbf{q}^\nu$	\mathbf{s}^μ	$\mathbf{1}$	$\mathbf{v}^\mu \mathbf{v}^\nu, \mathbf{s}^\mu \mathbf{s}^\nu$
$\mathcal{M}_N _{\text{SI}}$	$\vec{s} \cdot \vec{v}_\perp$	$\vec{s} \cdot \vec{v}_\perp$	1	$1, 1$

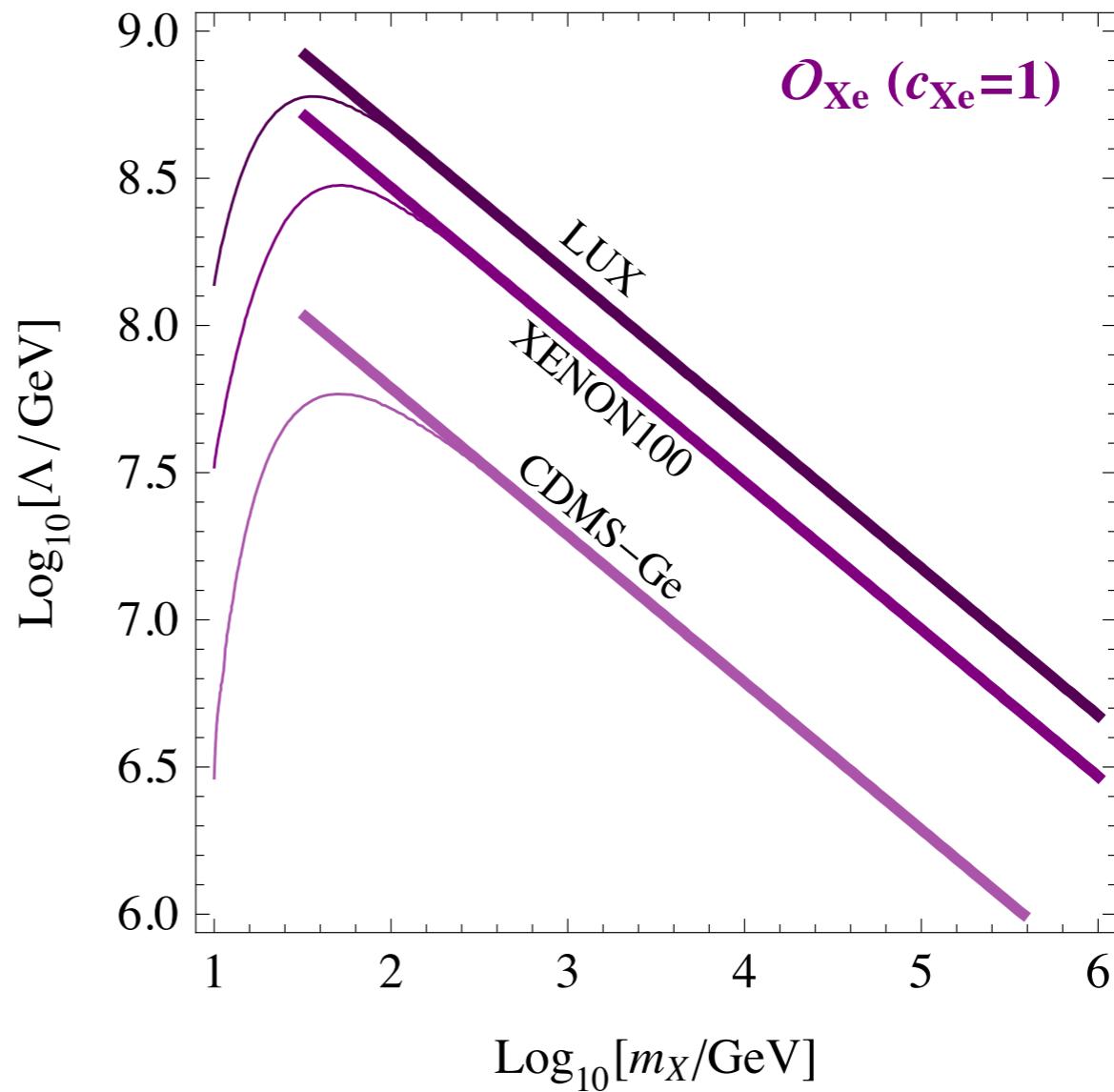
$\vec{v}^2, (\vec{s} \cdot \vec{v})^2, \dots$ **subleading** (advantage of relativistic formulation)

d=2 (dipole)

complex WIMPs with spin & generic CPV

$$\frac{c_{Xe}}{\Lambda} \bar{X} i\sigma^{\mu\nu} \gamma^5 X e F_{\mu\nu}$$

>1000 TeV!!!



$$c_{Xe} \sim \frac{g_*^2}{16\pi^2} \Rightarrow m_* \gtrsim 100 \text{ TeV } g_*^{4/3}$$

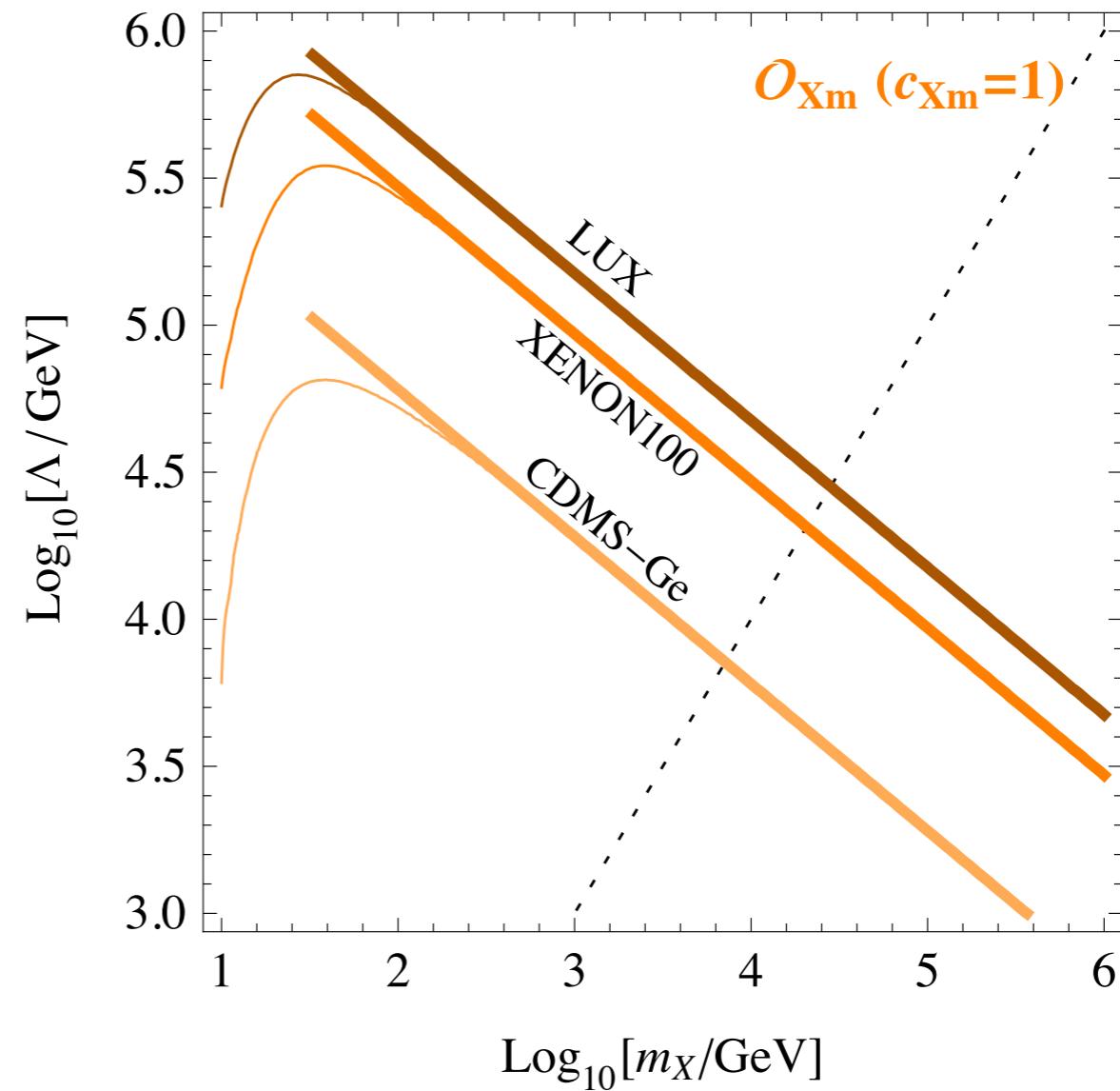
d=2 (dipole)

complex WIMPs with spin & ~~generic CPV~~

non-generic CPV is welcome (n-EDM)

$$\frac{c_{Xm}}{\Lambda} \bar{X} \sigma^{\mu\nu} X e F_{\mu\nu}$$

**>30 TeV
(>130 TeV XENON1T)**

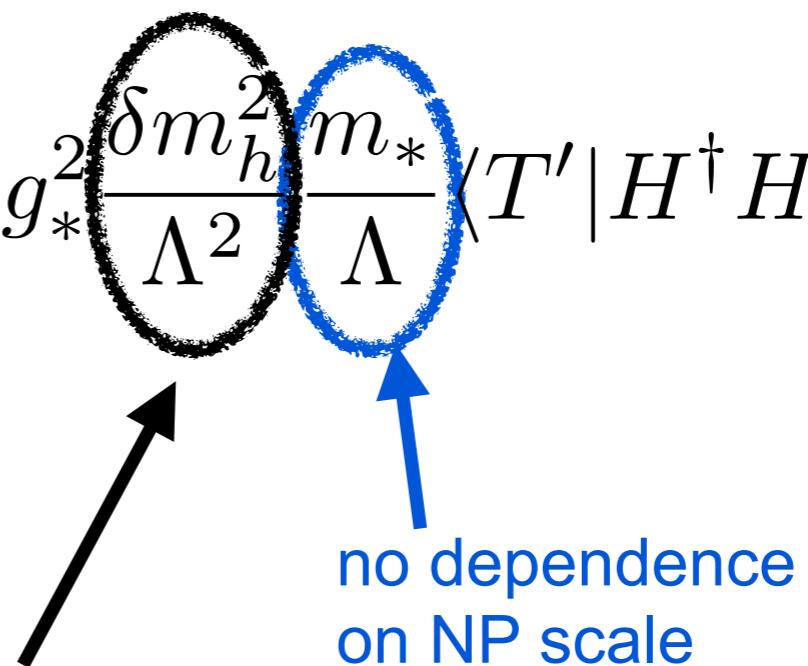


$$c_{Xm} \sim \frac{g_*^2}{16\pi^2} \Rightarrow m_* \gtrsim 4.5 \text{ TeV } g_*^{4/3} \text{ XENON1T}$$

d=2 (Higgs)

All WIMPs, Ex: fermion

$$\frac{g_*^2}{\Lambda} \frac{\delta m_h^2}{\Lambda^2} \bar{X} X H^\dagger H \implies \mathcal{M} \sim g_*^2 \frac{\delta m_h^2}{\Lambda^2} \frac{m_*}{\Lambda} \langle T' | H^\dagger H | T \rangle$$



- suppressed in PNGB Higgs typically $O(0.1-0.01)$
- enhanced by a light dilaton

$$m_* \gtrsim (0.5 - 2) \text{ TeV } g_*^{4/3}$$

XENONIT

d=3

complex WIMPs of any spin:

DD complementary to EW precision tests!

$$\frac{m_X}{\Lambda^2} X^\dagger v_\mu X e \partial_\nu B^{\mu\nu}$$



>2 TeV (LUX)
>5 TeV (XENON1T)

$$\frac{m_X}{f^2} X^\dagger v_\mu X H^\dagger i \vec{D}^\mu H$$



“T-parameter”

$$\frac{v^2}{f^2} \lesssim 10^{-3} \text{ (LUX)}$$

“Un-naturalness” at colliders

$$m_* \sim (3 - 5)g_*^2 \text{ TeV} \longrightarrow \frac{v^2}{f^2} \sim 5 \times 10^{-3} \frac{1}{g_*^2}$$

- * **no custodial SU(2)** -- truly accidental
more models, ex. minimal SU(3)/SU(2)xU(1)

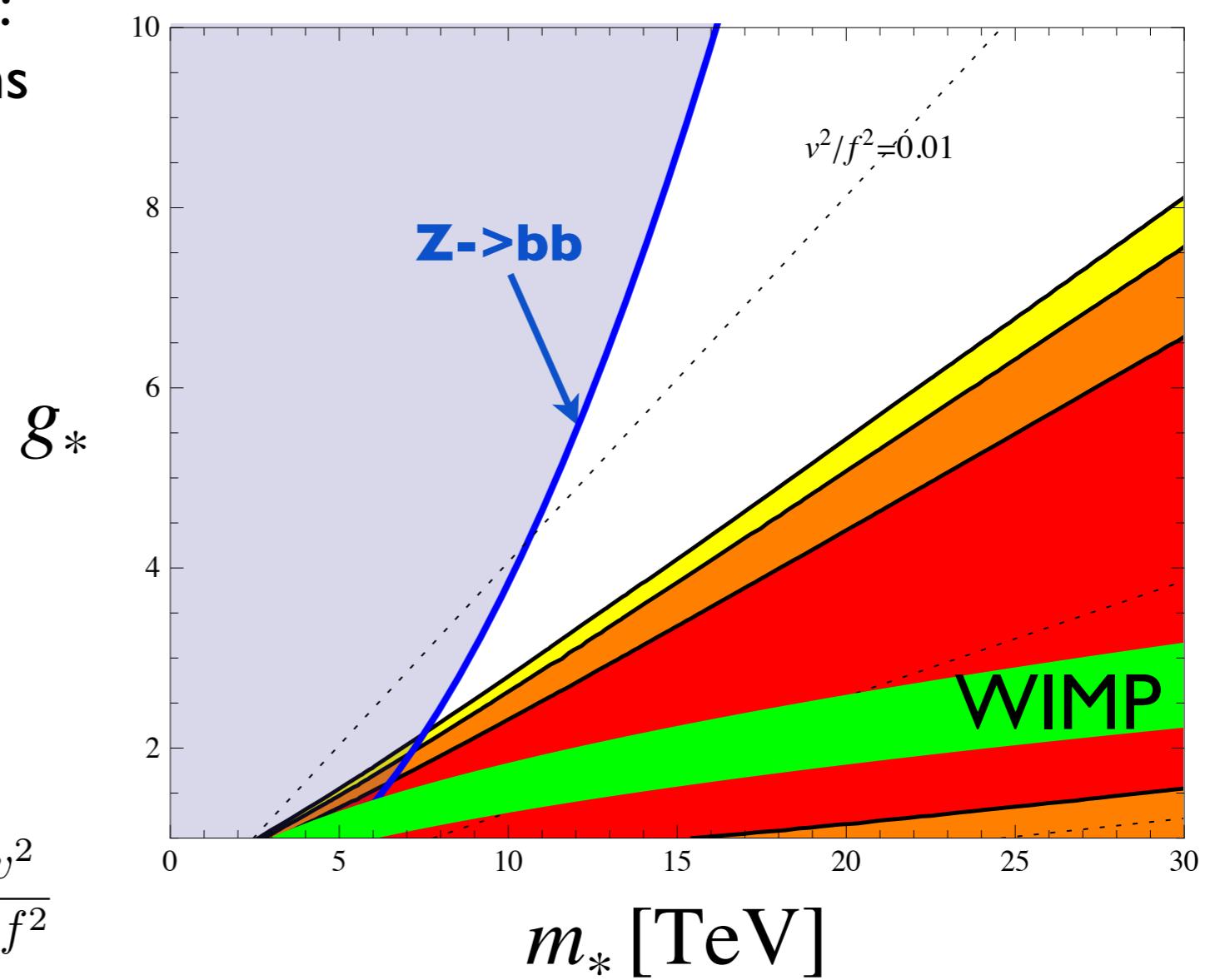
- * **no protection for Z-->bb:**
generic embedding of SM fermions

- * **Future linear collider?**

Higgs couplings >0.001

Abramowicz et al. (2013)

fit with $\alpha_{\text{em}} \Delta T_{\text{UV}} = +\frac{v^2}{4f^2}$



pp collider, resonant production

- Light dilaton? Davoudiasl et al. (2012)

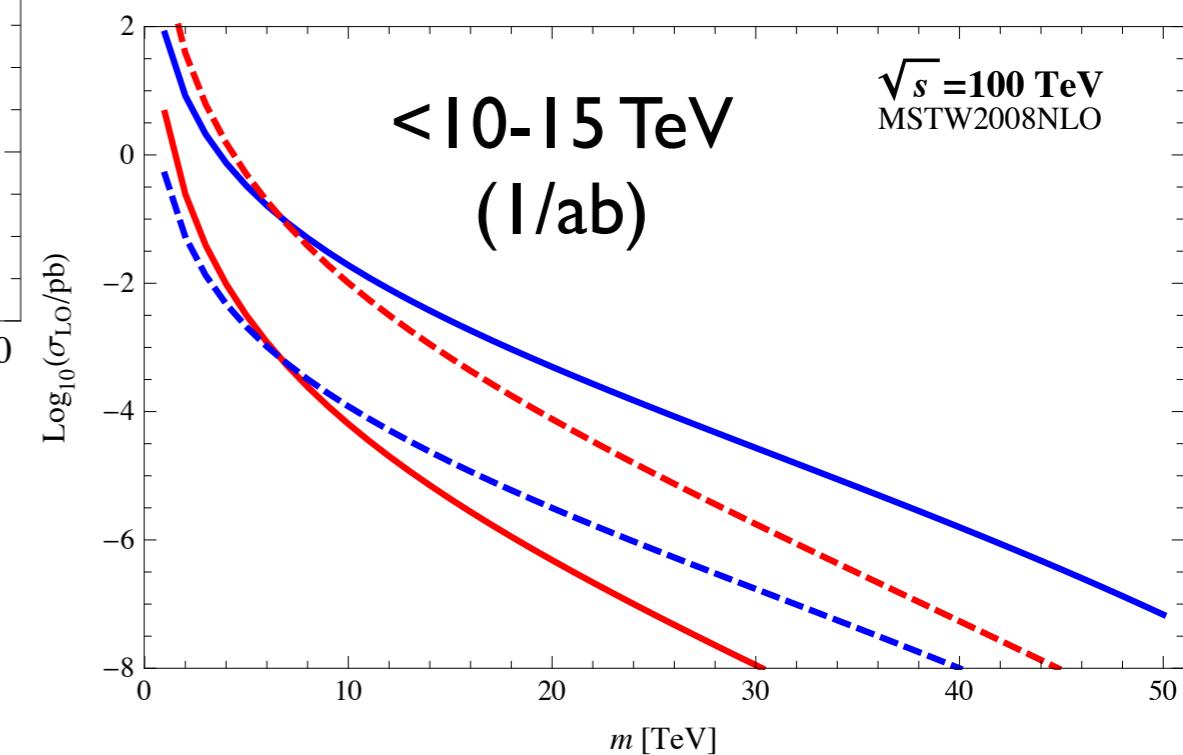
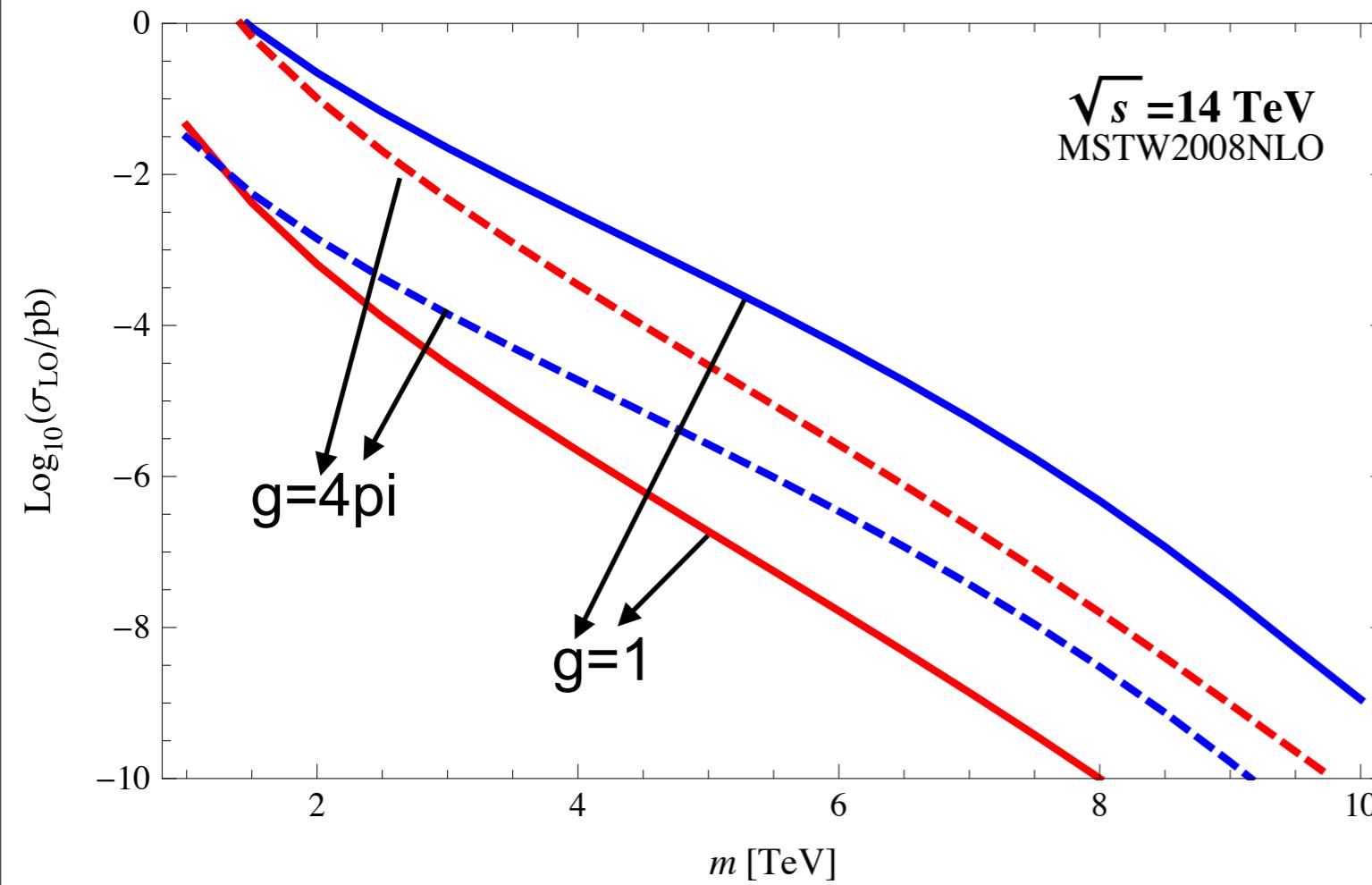
$$\frac{\alpha_s}{4\pi} GG \frac{\sigma}{f}$$

rough estimate:
 $m=f < 5 \text{ TeV}$, for 100 TeV and $1/\text{ab}$

- Light axions? $\frac{\alpha_s}{4\pi} F \tilde{F} \frac{a}{f}$ (non-minimal models)

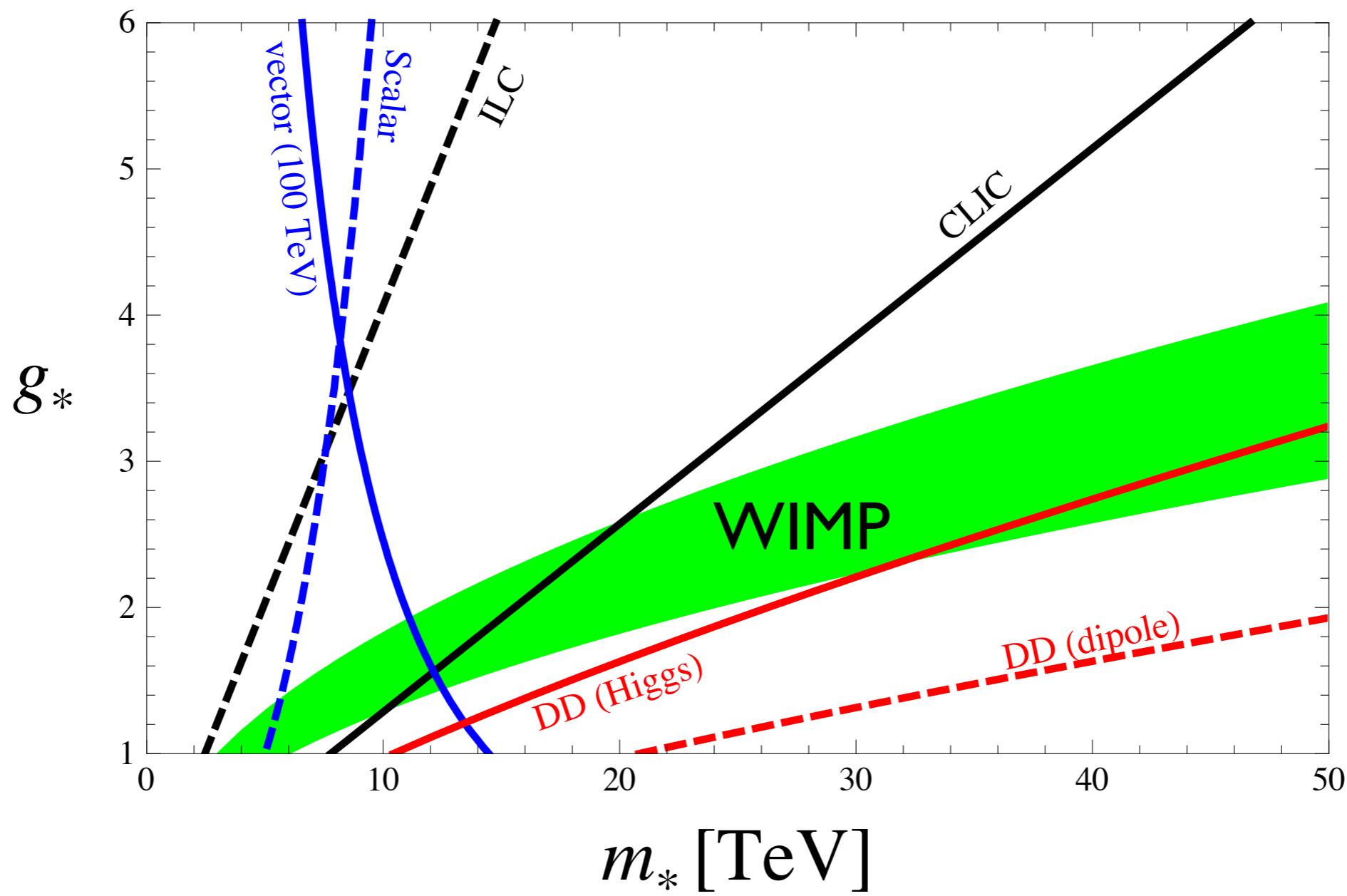
vector or scalar ?

$$\frac{g^2}{g_*} \bar{q} \gamma^\mu q \rho_\mu \quad \frac{\alpha_s}{4\pi} G G \frac{\sigma}{f}$$



vectors:

Dobrescu, Yu (2013), Pappadopulo et al. (2014), etc.



Conclusion

- * **Tuning can be natural:** profound connection Dark Matter \leftrightarrow Weak Scale
- * **Tuning can be predictive:** Dark matter is a key signature
Direct detection is a complementary (and very efficient) probe of un-naturalness
- * **Tuning can be attractive:**
pNGB Higgs \Rightarrow simple theory of Weak scale (& Higgs mass), Flavor, and DM

Back-up slides

EWPT

- S-parameter: $m > 3-5 \text{ TeV}$

$$\Delta S = 8\pi \frac{v^2}{m_\rho^2}$$

- T-parameter (model-dependent)

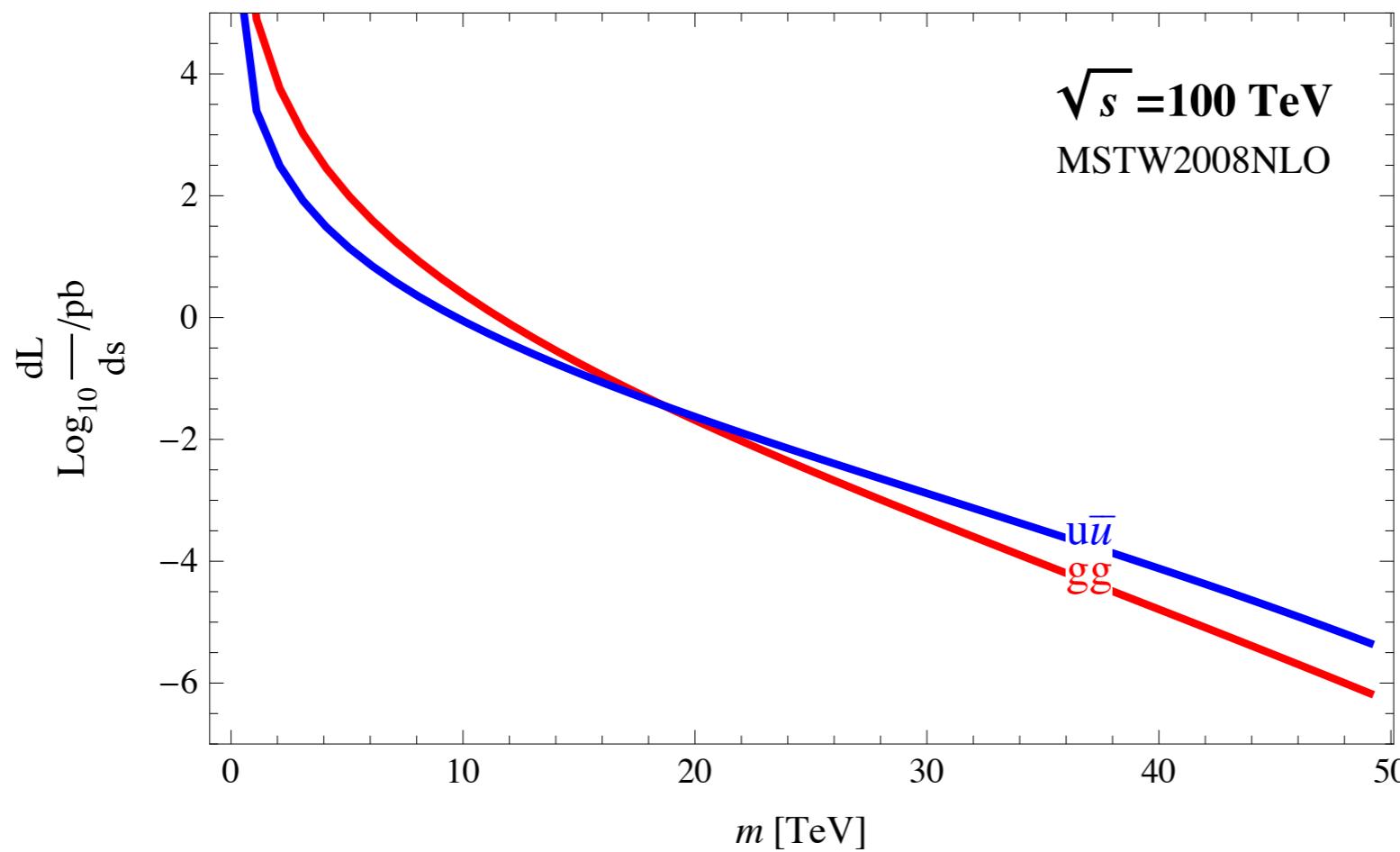
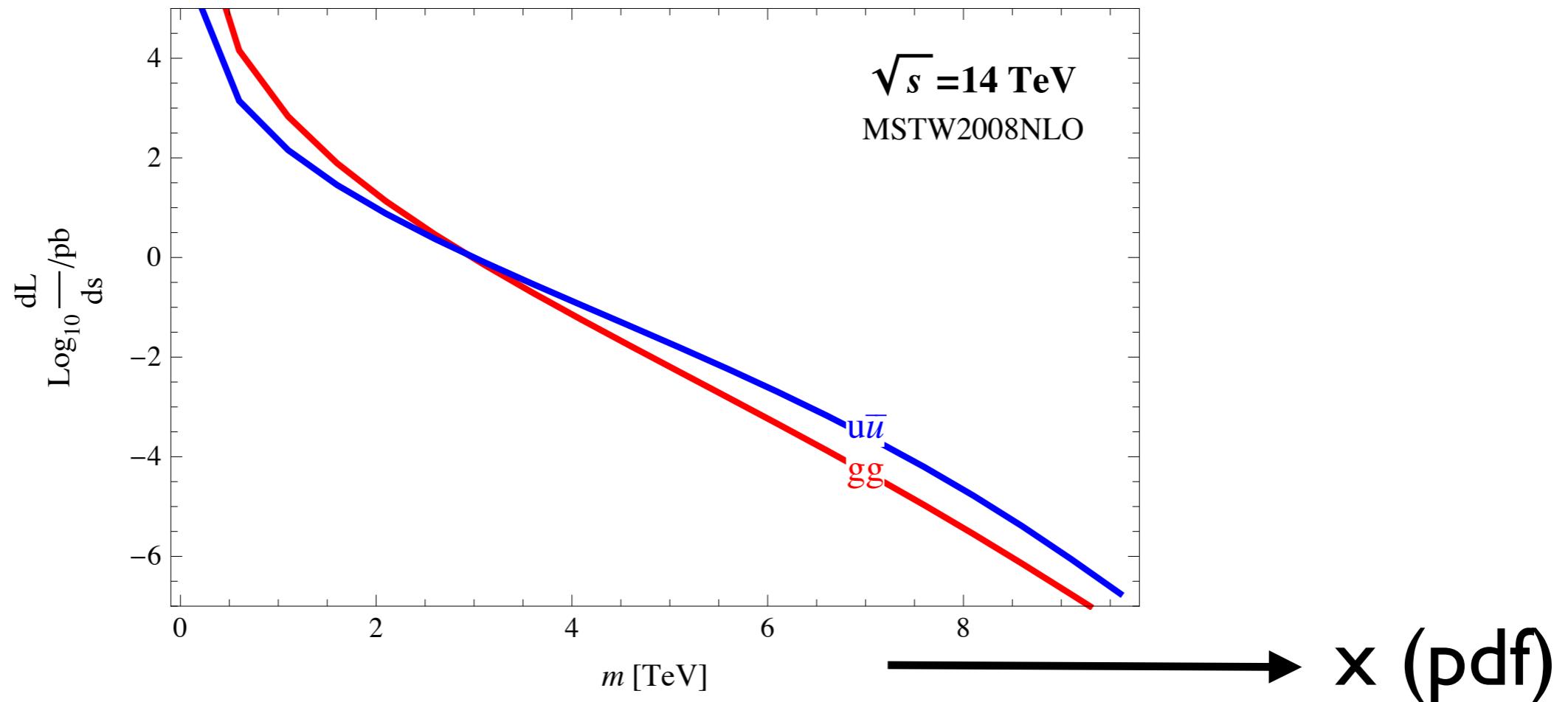
$$\Delta T = -\frac{3}{8\pi c_w^2} \frac{v^2}{f^2} \ln \frac{m_\rho}{m_h} + \Delta T_{\text{UV}}$$

125 GeV Higgs & Tuning

$$V = \frac{y_t^2 N_c}{16\pi^2} m_*^2 f^2 \left[a \sin^2 \frac{h}{f} + b \sin^4 \frac{h}{f} + \dots \right] \text{ plus gauge and/or exotic contributions}$$

total tuning $\sim \frac{\delta m_h^2}{m_h^2} \frac{\delta \lambda}{\lambda} \sim \frac{\delta a}{a} \frac{\delta b}{b} \sim \frac{m_*^2}{m_t^2} \left(\delta b \frac{g_*^2}{y_t^2} \right)$

~ 1
↑
 a b
↓
 v/f
↓
Higgs mass



Flavor and CP violation

$$m_* \sim (3 - 5)g_*^2 \text{ TeV}$$



- * **FV in the quark sector is under control for couplings >1-2**
- * **neutron EDM requires couplings >2-3 or non-generic CPV**
- * **mu--->e gamma and e-EDM, three options:**
 - 1) maximal coupling
 - 2) non-generic FV, CPV
 - 3) just luck at % level